

- ideal gas equation of state from kinetic theory

$$PV = Nk_B T \quad (1)$$

$$\frac{1}{2}m\overline{v^2} \equiv \frac{3}{2}k_B T \quad (2)$$

- a theorem (that relates susceptibilities)

$$f(x, y, z) = 0$$

$$\left. \frac{\partial x}{\partial y} \right|_z \left. \frac{\partial y}{\partial z} \right|_x = - \left. \frac{\partial x}{\partial z} \right|_y \quad (3)$$

- 1st law - conservation of energy

$$dU = \delta W + \delta Q$$

- reversible work $W_{AB} = - \int_A^B P dV$

- specific heat -

$$C_V = \frac{1}{M} \left. \frac{\delta Q}{dT} \right|_V = \frac{1}{M} \left. \frac{\partial U}{\partial T} \right|_V \quad (4)$$

$$C_P = \frac{1}{M} \left. \frac{\delta Q}{dT} \right|_P = \frac{1}{M} \left[\left. \frac{\partial U}{\partial T} \right|_P + P \left. \frac{\partial V}{\partial T} \right|_P \right] \quad (5)$$

- ideal gas - $U(T)$

$$U = \frac{3}{2} N k_B T \text{ - monoatomic}$$

$$U = \frac{5}{2} N k_B T \text{ - diatomic}$$

- $S = k_B \ln \Gamma$ - Boltzmann (microcanonical) entropy

$$T \equiv \left. \frac{\partial U}{\partial S} \right|_{P, V}$$

$$P \equiv \left. \frac{\partial U}{\partial V} \right|_{S, N}$$

- thermodynamic identity

$$dU = \left. \frac{\partial U}{\partial S} \right|_V dS + \left. \frac{\partial U}{\partial V} \right|_S dV = T dS - P dV \quad (6)$$

- for reversible process $dS = \frac{\delta Q}{T}$, $\delta W = -PdV$
- in an irreversible process

$$\Delta S > \frac{\delta Q}{T} \quad (7)$$

- in a cyclic transformation $\Delta S = 0$
irreversible

$$\sum_i \frac{Q_i}{T_i} < 0 \quad (8)$$

reversible

$$\sum_i \frac{Q_i}{T_i} = 0 \quad (9)$$

- cyclic Carnot engine - efficiency η

$$\eta = \frac{|W|}{|Q_h|}, \eta \leq 1 - \frac{T_c}{T_h} \quad (10)$$

- for reversible engine $\eta = 1 - \frac{T_c}{T_h}$
- cyclic Carnot refrigerator - efficiency η'

$$\eta' = \frac{|Q_c|}{|W|}, \eta' \leq \frac{T_c}{T_h - T_c} \quad (11)$$

- Thermodynamic Functions \rightarrow Maxwell relations
- internal energy $U(S, V), dU = TdS - PdV$
- "Helmholtz" free energy $F(T, V) = U - TS, dF = -SdT - PdV$
- enthalpy $H(P, S) = U + PV, dH = TdS + VdP$
- "Gibbs" free energy $G(P, T) = U + PV - TS, dG = VdP - SdT$
- Clausius-Clapeyron relation $\frac{dP}{dT} = \frac{L}{T\Delta V}$
- Canonical Ensemble
probability to find system at state n of energy ϵ_n

$$p_n = \frac{e^{-\frac{\epsilon_n}{k_B T}}}{Z} \quad (12)$$

- partition function $Z = \sum_n e^{-\frac{\epsilon_n}{k_B T}}$

$$\frac{p_n}{p_m} = \frac{e^{-\frac{\epsilon_n}{k_B T}}}{e^{-\frac{\epsilon_m}{k_B T}}} = e^{-\frac{\epsilon_n - \epsilon_m}{k_B T}} \quad (13)$$

$$U = \sum_n p_n \epsilon_n \quad (14)$$

$$S = -k_B \sum_n p_n \ln p_n \quad (15)$$

- free energy $F = U - TS = -k_B T \ln Z$
- free energy (and thus Z) generates all thermodynamic quantities

$$S = -\left. \frac{\partial F}{\partial T} \right|_V \quad (16)$$

$$U = -T^2 \left. \frac{\partial}{\partial T} \left(\frac{F}{T} \right) \right|_V \quad (17)$$

$$C_V = -T^2 \left. \frac{\partial^2 F}{\partial T^2} \right|_V \quad (18)$$

- Boltzmann distribution- (p,q)

$$dw = \frac{e^{-\epsilon(p,q)/k_B T}}{Z} dpdq$$

$$Z = \frac{1}{h} \int dpdq e^{-\epsilon(p,q)/k_B T}$$

$$\bar{f} = \frac{\int dpdq f(p,q) e^{-\beta \epsilon(p,q)}}{\int dpdq e^{-\beta \epsilon(p,q)}}$$

- Grand Canonical Ensemble
probability to find system at state n_N of N particles and energy E_{n_N}

$$p_{n_N} = A e^{-\beta(E_{n_N} - \mu N)} \quad (19)$$

- Landau potential

$$\Omega = -k_B T \ln \sum_N \sum_{n_N} e^{-\beta(E_{n_N} - \mu N)} \quad (20)$$

- independent quantum particles
Bose-Einstein statistics

$$n_\nu = \frac{1}{e^{\beta(\epsilon_\nu - \mu)} - 1} \quad (21)$$

Fermi-Dirac statistics

$$n_\nu = \frac{1}{1 + e^{\beta(\epsilon_\nu - \mu)}} \quad (22)$$

- translationally invariant system

$$\sum_k \rightarrow \frac{L}{2\pi} \int dk \quad (23)$$